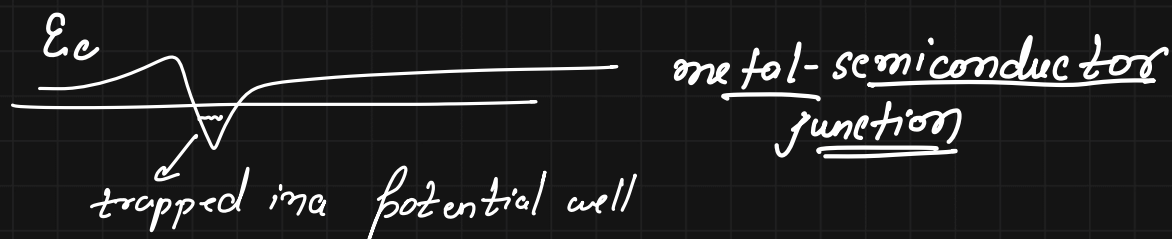
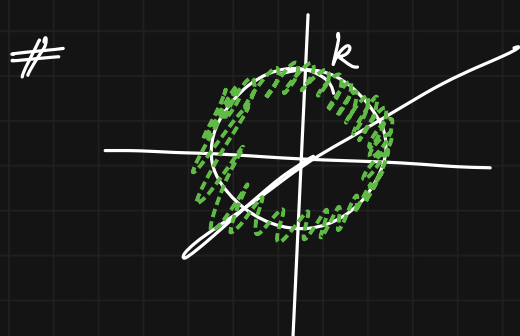


Recap

2D system $\rightarrow$ not only where topo is relevant



drift $\rightarrow$ drude picture $\rightarrow$ elastic collisions, momentum gets scrambled

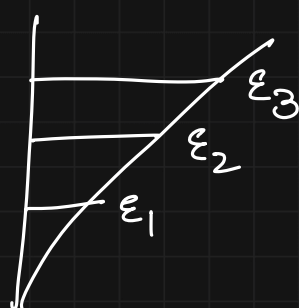


$$\left(\frac{dp}{dt}\right)_{\text{scattering}} = \left(\frac{dp}{dt}\right)_{\text{field}}$$

$$v_d = e \frac{2m}{m} E \quad \text{or} \quad \mu = \frac{e 2m}{m}$$

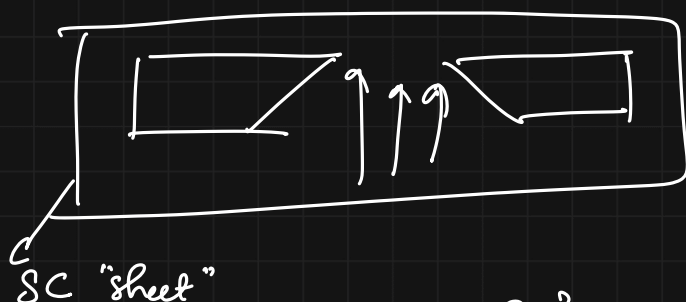
concept of subbands $\rightarrow$ for confinement along z-direction

$$\psi(z) = \phi_n(z) e^{ik_x x} e^{ik_y y}$$



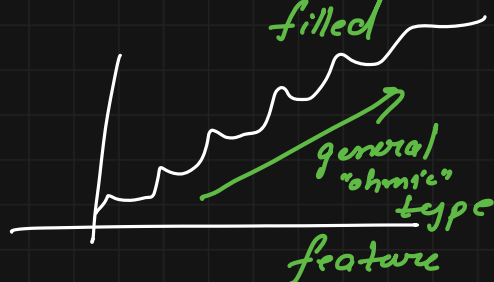
$$E = E_c + E_n + \frac{\hbar^2}{2m} (k_x^2 + k_y^2)$$

PRL paper $\rightarrow$ quantized conductance



$$\sigma = \frac{2e^2}{h} \eta$$

Gate voltage $\rightarrow$ # of states that are filled



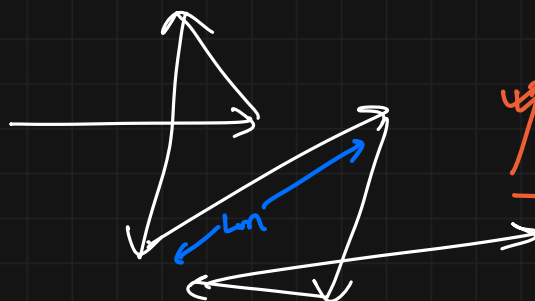
Question 2

diffusive model works? or do we need something extra?

remark

length scales

momentum scattering length.



$$l_m \sim v_f \tau_m$$

"This is due to ELASTIC scattering" & not INELASTIC scattering

Problem:-

$$\# \text{ states} \rightarrow \frac{1}{4\pi^2} \times 2 \times \underline{d^2k} \quad \mathcal{E} = \frac{\hbar^2 k^2}{2m}$$

$$= \frac{1}{2\pi^2} \cdot k dk (2\pi)$$

$$= \frac{1}{\pi} k dk$$

$$d\mathcal{E} = \frac{\hbar^2}{m} k dk$$

$$= \left(\frac{1}{\pi} \right) \underline{d\mathcal{E}} \frac{m}{\hbar^2} \Rightarrow \frac{m}{\pi \hbar^2} = g(\mathcal{E})$$

$$\int \frac{1}{\pi} \times (1) \times d\mathcal{E} = \frac{m}{\hbar^2} \frac{\mathcal{E}_F}{\pi} = \eta \quad \begin{array}{l} \downarrow \\ \text{indpt of} \\ \text{energy} \end{array}$$

$$\frac{m}{\hbar^2} \times \frac{1}{\pi} \frac{\hbar^2 k_F^2}{2m} = \eta$$

$$k_F = \sqrt{2\pi\eta}$$

periodic BCs

$$\boxed{\frac{d\eta}{d\mathcal{E}} = \frac{m}{\pi \hbar^2}}$$

$$m v_f = \hbar k_F \Rightarrow v_f = \frac{\hbar k_F}{m} = \frac{\hbar}{m} \sqrt{2\pi\eta}$$

$$\boxed{l_m = v_f \tau_m \rightarrow \text{momentum scattering}}$$

1 more length scale

$$\lambda_f = \frac{1}{k_F} \rightarrow \lambda_f = \text{de-broglie wavelength of electrons}$$

$$\lambda_f \sim (n^{-1/2})$$

$$k_f \sim (n)^{1/d}$$

$$\lambda_f \sim (n)^{-1/d}$$

Problem

What is the typical λ_f for electrons in typical metals?

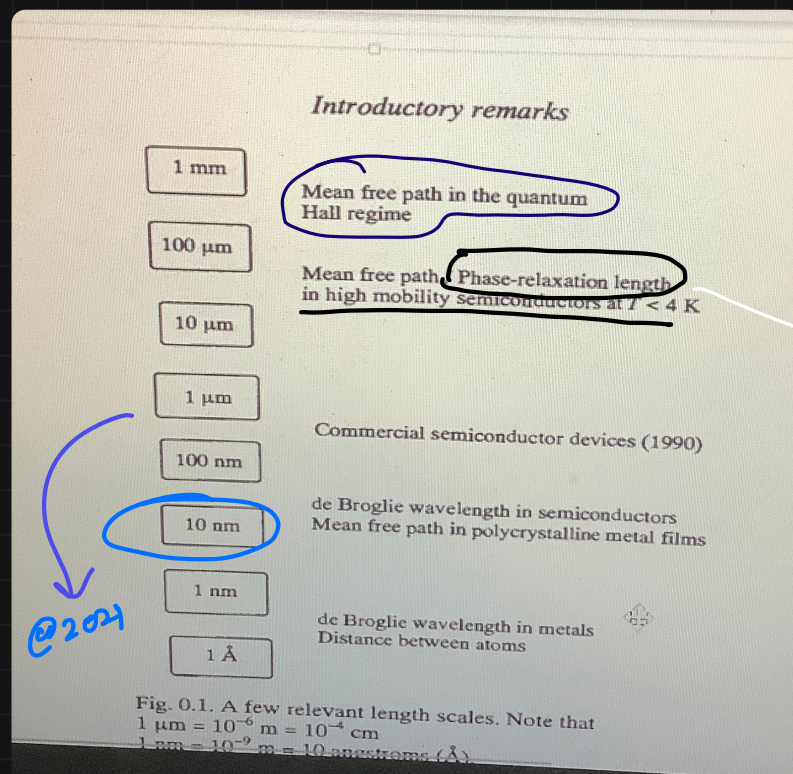
Ans:- $\lambda_f \sim \text{\AA}$ (for typical metals) $n \sim 10^{22} \text{ cm}^{-3}$

atomic spacing

$$\lambda_f \sim (n)^{-1/d}$$

$\lambda_{f \text{ semicond}} \sim \text{larger than metals}$

$\rightarrow$ essentially small n



$\rightarrow$ S. datta

gives diff regions of transport

@200

Phase relaxatⁿ

If source of collision is inelastic $\rightarrow$ "phase" of e gets scrambled

Phase coherence time scale:-

drop an e^- into a solid $\rightarrow$ gets scattered due to "inelastic collisions"

$\swarrow$ (e^- , phonons, magnetic)

Quantum coherence of phase is lost.

(decoherence/dephasing)

elastic scattering $\rightarrow$ no decoherence

$\rightarrow$ might add a const. phase

on some timescale $\tau_\phi \rightarrow e^-$ becomes uncorrelated from its initial phase.

$\swarrow$ interference effects wash out.

Dist e^- moves in this time in a diffusive system

$$l_\phi = (D \tau_\phi)^{1/2} = \text{"Coherence length"}$$

Length scales

As inelastic processes freeze out, τ_ϕ is expected to diverge, with power law exponent set by dimensionality and inelastic mechanism:

electron phonon: $\tau_\phi \sim T^{-3}$

electron-electron: $\tau_\phi \sim T^{2(d-4)}$

Phase coherence length $l_\phi = (D \tau_\phi)^{1/2} \uparrow$ as $T \downarrow$

Mean free path $l_e = v_F \tau$ dominated by elastic scattering (impurities, roughness) and is stays constant

mesoscopic physics regime

Sample channel length L

$L < l_e$, ballistic,

$L < l_\phi$, phase coherent,

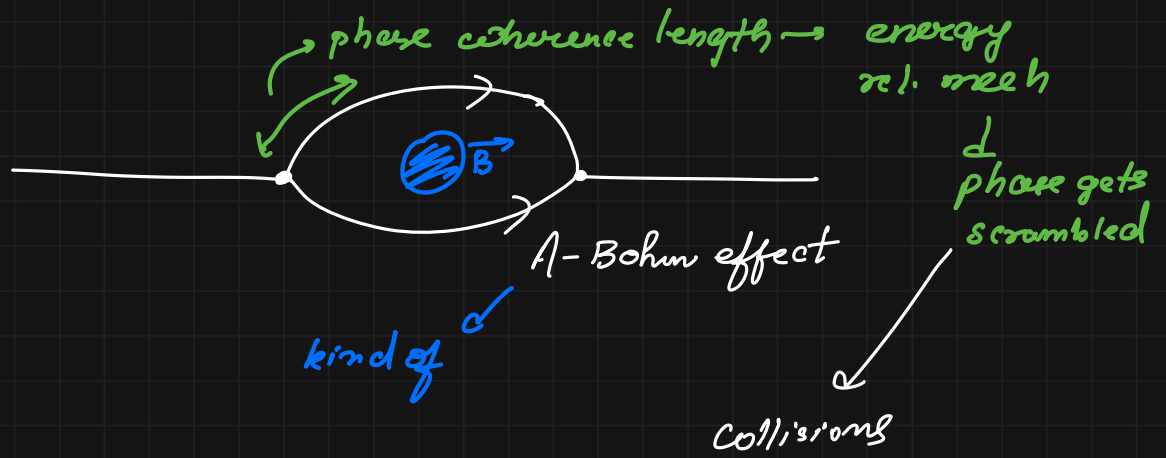
$L > l_e$, diffusive

$L > l_\phi$, incoherent

$l_\phi > l_e$ at low temperatures, and lots of interesting phenomena observed when $l_e < L < l_\phi$

$\rightarrow$ weak-anti-localization

expt $\rightarrow$ double slit

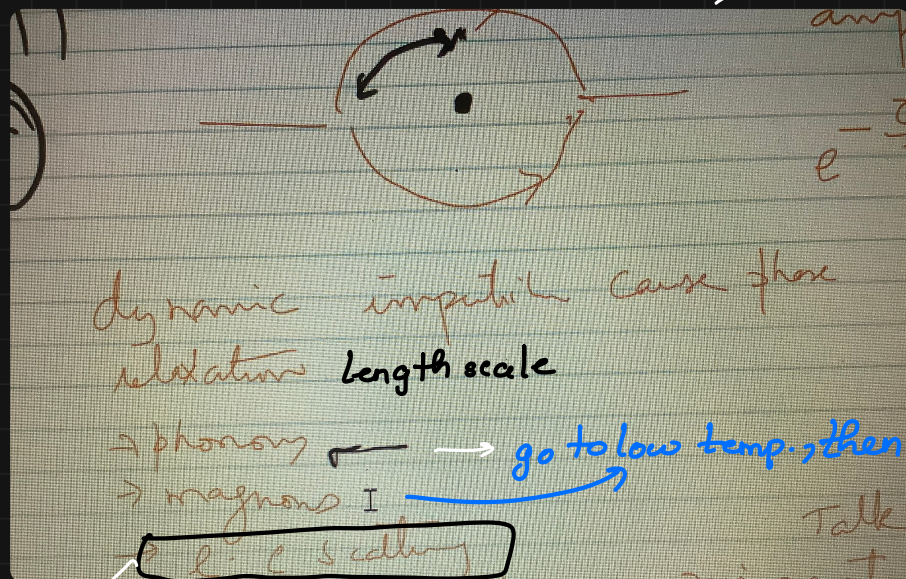
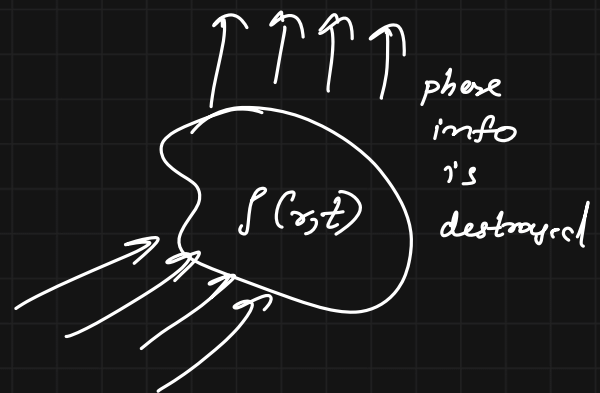
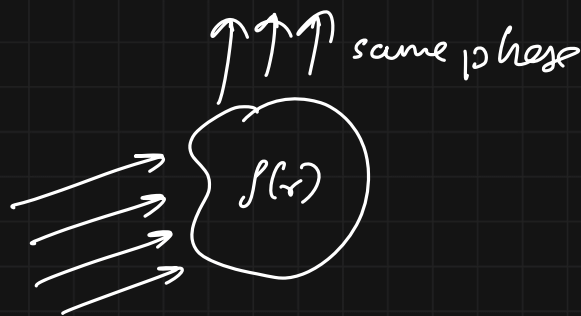


B16
picture

All collisions that have fixed target, then the phase is same in & out.

$\hookrightarrow$ "no time dependence"

$\hookrightarrow$ t dependence $\Rightarrow$



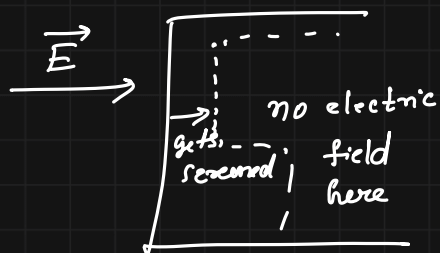
$\hookrightarrow$ at low temp, it's still there.

$\hookrightarrow$ even in non-correlated systems, it's there.

$\downarrow$
"metals"

1 more
lengthscale

Screening length λ_s



diff models
 $\hookrightarrow$ Thomas-Fermi model
 $\lambda_{TF} \sim (\lambda_F)^{\text{power}} \times \text{const.}$
 $\sim \text{few } \text{\AA} \text{ in metals}$

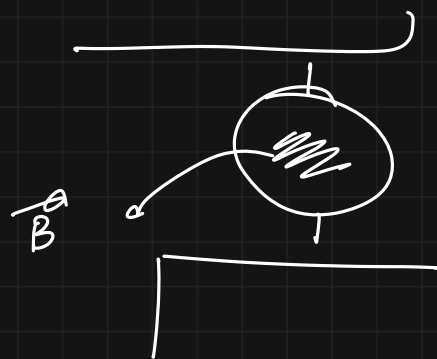
in SC the lengthscale can be large

e.g. for transistors $\rightarrow$ have a large screening "length"

$\lambda_F < \text{atomic spacing} \rightarrow \text{particle in a box}$

Paper
measuring
phase
coherence
length

Mesoscopic "physics"



G vs B_{ext}

$\downarrow$
faster beating oscillations

$\hookrightarrow$ a resetting of phase & oscillation

$\text{dia} < \text{en. rel. length} \Rightarrow \text{oscillations coarsen away}$

Φ_0 component $\rightarrow T = \text{knob for phonons}$

amplitude of oscillations coarsens away

Control inelastic collisions

→ Datta's Book → phase coherence length

A2M

⇒ L. length → not observed for short magnetic field.

↓
always about the numbers

